

Testing Composite Higgs models on the lattice

Enrico Rinaldi

*Physics Division,
Lawrence Livermore National Laboratory,
Livermore, CA*

Lattice Strong Dynamics collaboration

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Lattice **S**trong **D**ynamics Collaboration



James Osborn



Evan Berkowitz
Enrico Rinaldi
Chris Schroeder
Pavlos Vranas



Rich Brower
Michael Cheng
Claudio Rebbi
Oliver Witzel
Evan Weinberg



Joe Kiskis



Ethan Neil



David Schaich



Ethan Neil
Sergey Syritsyn



Tom Appelquist
George Fleming
Gennady Voronov



Meifeng Lin



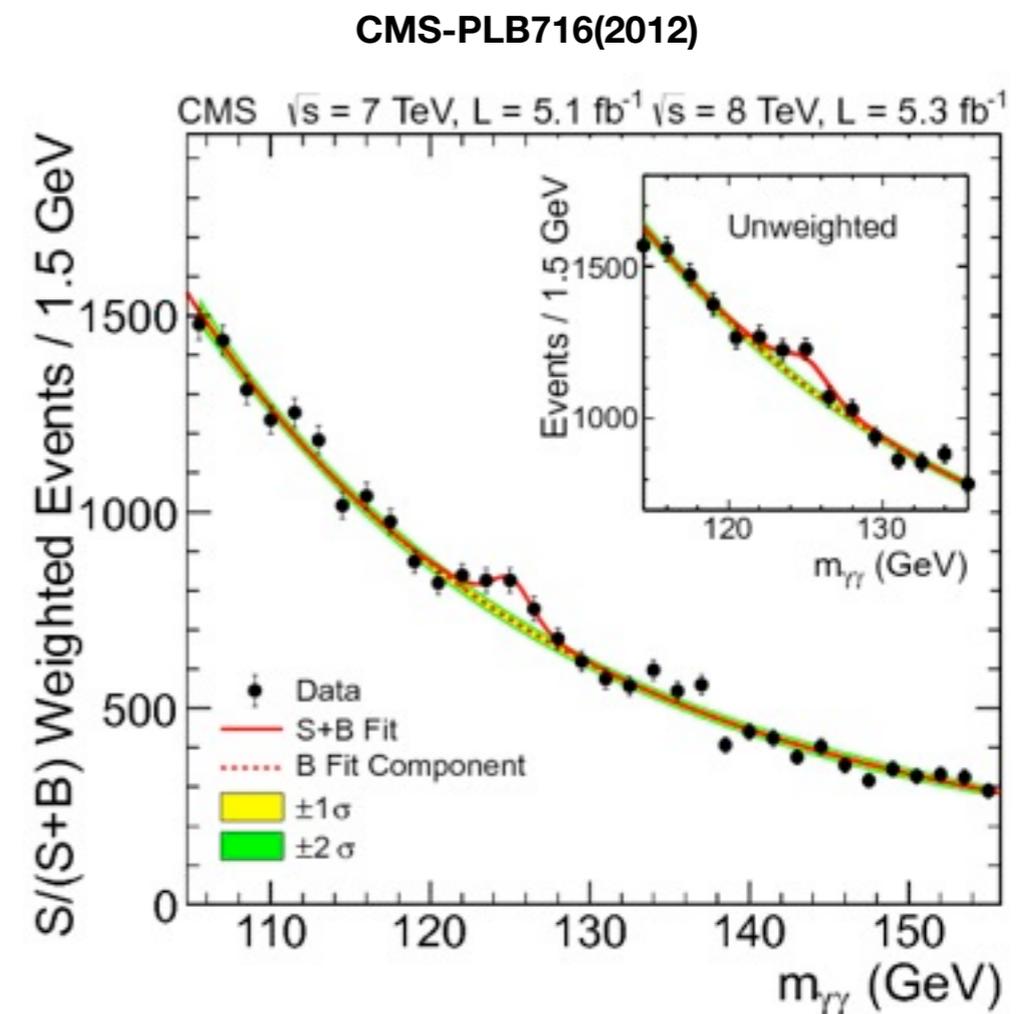
Graham Kribs



Mike Buchoff

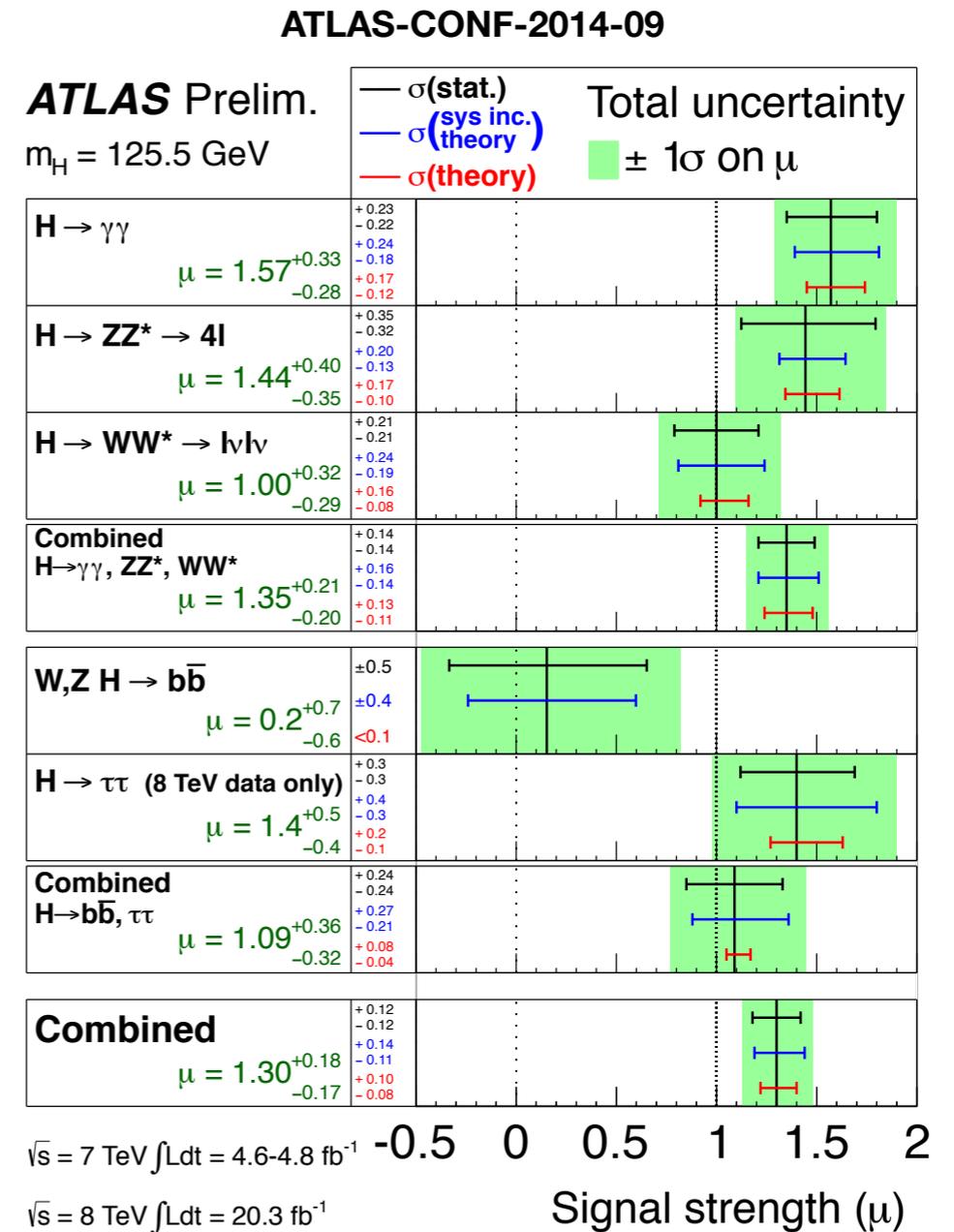
The SM Higgs boson

- Higgs boson discovery in Run I at LHC
↳ $m_H = 125.5 \text{ GeV}$
- Higgs boson decays to ↳ **Gauge bosons** and **leptons+quarks**
- Signal strength measurements ↳ coupling constant measurements
- Run II at LHC ↳ **deviations from SM couplings? NP resonances above 1TeV?**
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- fine tuning of the EW scale
- triviality of the SM Higgs quartic coupling
- SM vacuum instability
- flavor problem
- neutrino masses
- dark matter candidates

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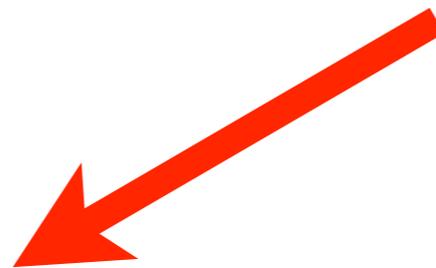
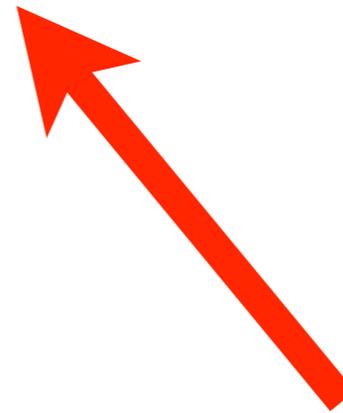
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compositeness of the Higgs boson

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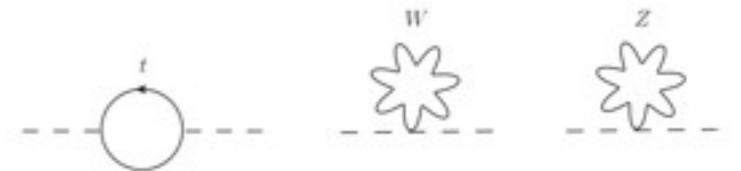
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Composite Higgs models

- the Higgs boson is not an **elementary** scalar particle \rightarrow **composite** bound state of new strong dynamics
- **Technicolor** Higgs:
 - ✓ the new sector breaks the EW symmetry through a technifermion condensate
 - ✓ the Higgs is identified with the lightest scalar excitation of the condensate
 - ✓ can be light due to interactions with SM particles (obtained from ETC dynamics)

[Foadi et al., PRD87(095001)] [Di Chiara et al., arxiv:1405.7154]



- **Walking Technicolor** Higgs:
 - ✓ walking coupling and large anomalous mass dimension $\gamma \approx 1$
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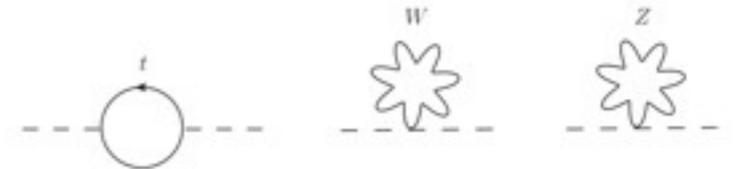
[Yamawaki et al., PRL56(1986)] [Bando et al., PLB178(1986)]

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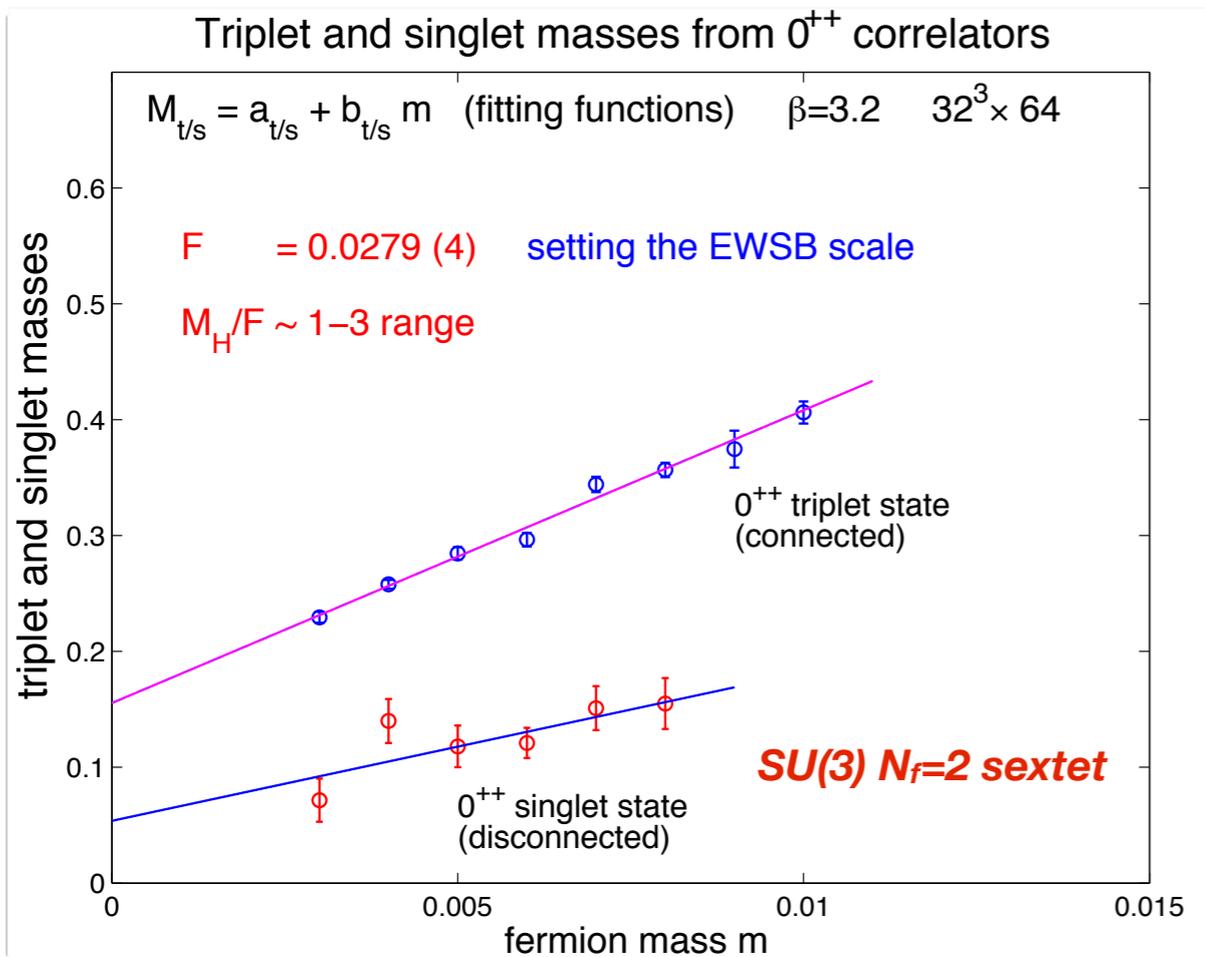
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Lattice simulations

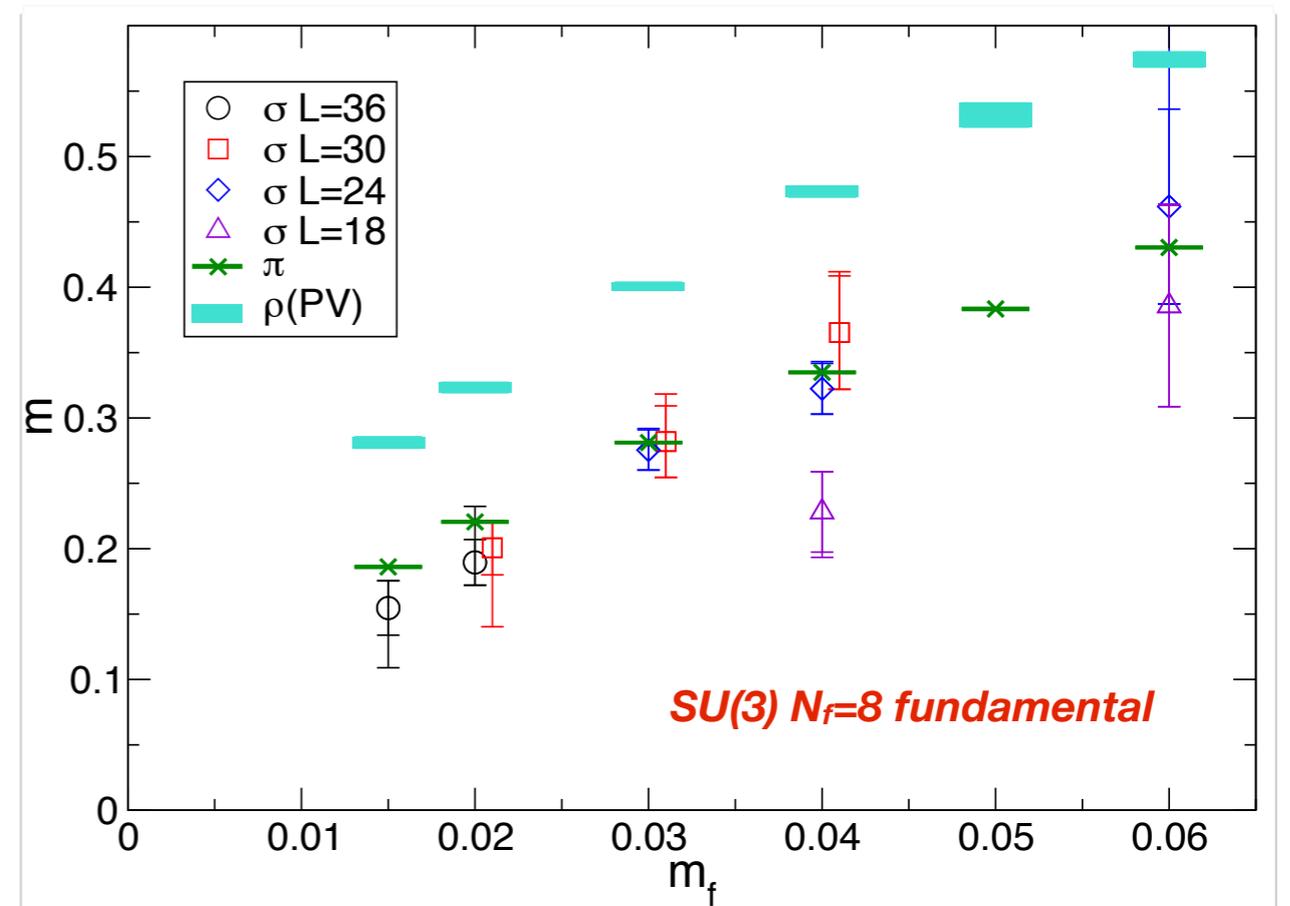
Composite Higgs models

LH Coll. PoS(2013)062



[see talk by R.Wong 2C for updates]

LatKMI Coll. PRD89(111502)R



[see talk by K.-i.Nagai 9C for updates]

Goldstone Higgs dynamics

- a composite Higgs can emerge as a **scalar** pseudo-NGB from breaking of a global symmetry in a new strong sector at the TeV scale
- phenomenologically interesting and well studied models are based on the breaking pattern **$SU(4) \rightarrow Sp(4)$** (equivalently $SO(6) \rightarrow SO(5)$) [Galloway et al., JHEP10(2010)89]
- a realization of this framework has been studied recently, showing how $SU(4)$ can break to $Sp(4)$ using a **$SU(2)$ gauge theory with 2 flavors in the fundamental representation** [Cacciapaglia&Sannino, JHEP04(2014)111]
 - ✓ $SU(4)/Sp(4)$ coset gives 5 NGBs \rightarrow 3 pseudoscalars and 2 scalars
 - ✓ different choices of the quark condensate can be used when embedding the strong sector with the EW sector
 - ✓ interplay between the **2 scalars NGBs** and the **lightest excitation of the condensate** depending on the vacuum alignment give different scenarios [talk by A.Hietanen 2C]

Higgs and dark matter candidates

- interest in this special theory is not limited to Higgs compositeness:
 - models for **composite DM** have been recently studied based on this framework, without connection with EW symmetry breaking [[Buckley&Neil, PRD87\(043510\)](#)]
 - a model for **dark nucleosynthesis** based on the same strong sector now exists [[Detmold et al., arxiv:1406.2276,arxiv:1406.4116](#)] [talk by W.Detmold 8C]
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study the isosinglet scalar channel on the lattice

Lattice simulations with FUEL

- FUEL (Framework for Unified Evolution of Lattices) it's a lightweight and flexible wrapper for level 3 USQCD libraries with focus on lattice generation for BSM theories [\[https://github.com/jcosborn/qhmc\]](https://github.com/jcosborn/qhmc) [See J. Osborn's talk in 1F]
- capabilities to do arbitrary number of **colors**, **dimensions** and **flavors**
- Staggered fermion formulations, as well as Wilson ones are included
- use HMC to generate $O(10^3)$ trajectories for **SU(2) with Wilson plaquette gauge action and 2 Wilson fermions**
- one coupling $\beta=2.2$, one volume $32^3 \times 64$, six bare fermion masses $m_0=\{-0.68,-0.70,-0.72,-0.735,-0.75,-0.755\}$

Glueball spectroscopy: operators

- eigenstates of the Hamiltonian are classified according to the irreducible representations of the cubic group

$$\{A_1(1), A_2(1), E(2), T_1(3), T_2(3)\}$$

- suitable gauge-invariant operators must be constructed that respect the symmetries

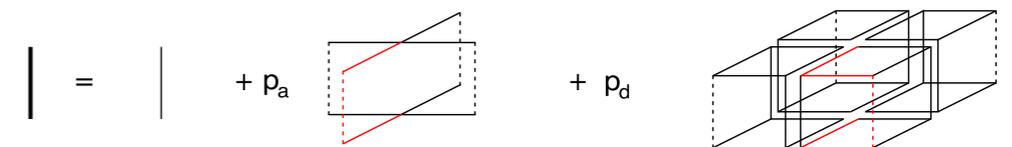
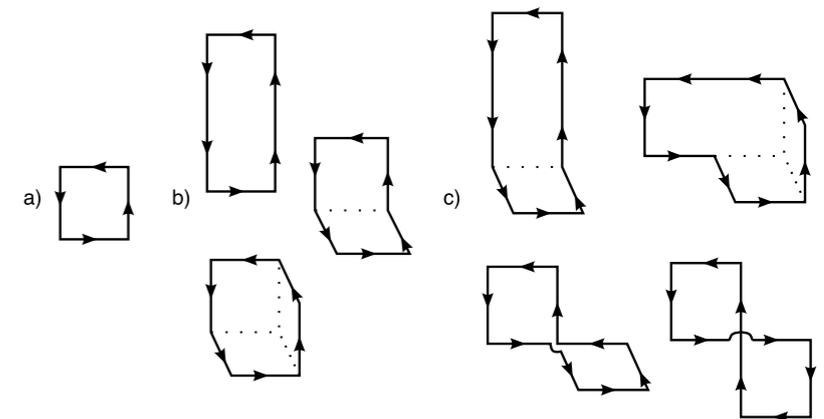
$$\mathcal{O}_G(t) = \frac{1}{L^3} \sum_{x \in L^3} \text{Tr} \left(\prod_{l \in \mathcal{W}(x)} U_l \right) \quad \mathcal{O}_G^{(R)}(t) = \sum_{\alpha=1}^{24} a_\alpha^{(R)} \mathcal{R}_\alpha [\mathcal{O}_G(t)]$$

- vacuum contributions must be subtracted in the scalar case

$$\mathcal{O}^{(A_1)}(t) - \langle 0 | \mathcal{O}^{(A_1)} | 0 \rangle$$

- improved operators are obtained by blocking and smearing algorithms

J	A_1	A_2	E	T_1	T_2
0	1	0	0	0	0
1	0	0	0	1	0
2	0	0	1	0	1
3	0	1	0	1	1
4	1	0	1	1	1



Lucini, Rago, ER
JHEP08(2010)

Glueball spectroscopy: variational analysis

- basis of operators \rightarrow

$$\{\mathcal{O}_1(t), \dots, \mathcal{O}_n(t)\}$$

- matrix of correlators \rightarrow

$$C_{ij}(t) = \sum_{\tau} \langle 0 | \mathcal{O}_i^\dagger(\tau + t) \mathcal{O}_j(\tau) | 0 \rangle$$

- generalised eigenvalue problem \rightarrow

$$C_{ij}(t) v_j^\alpha = \lambda^\alpha v_i^\alpha$$

- ground state correlator fit ($\alpha=0$) \rightarrow

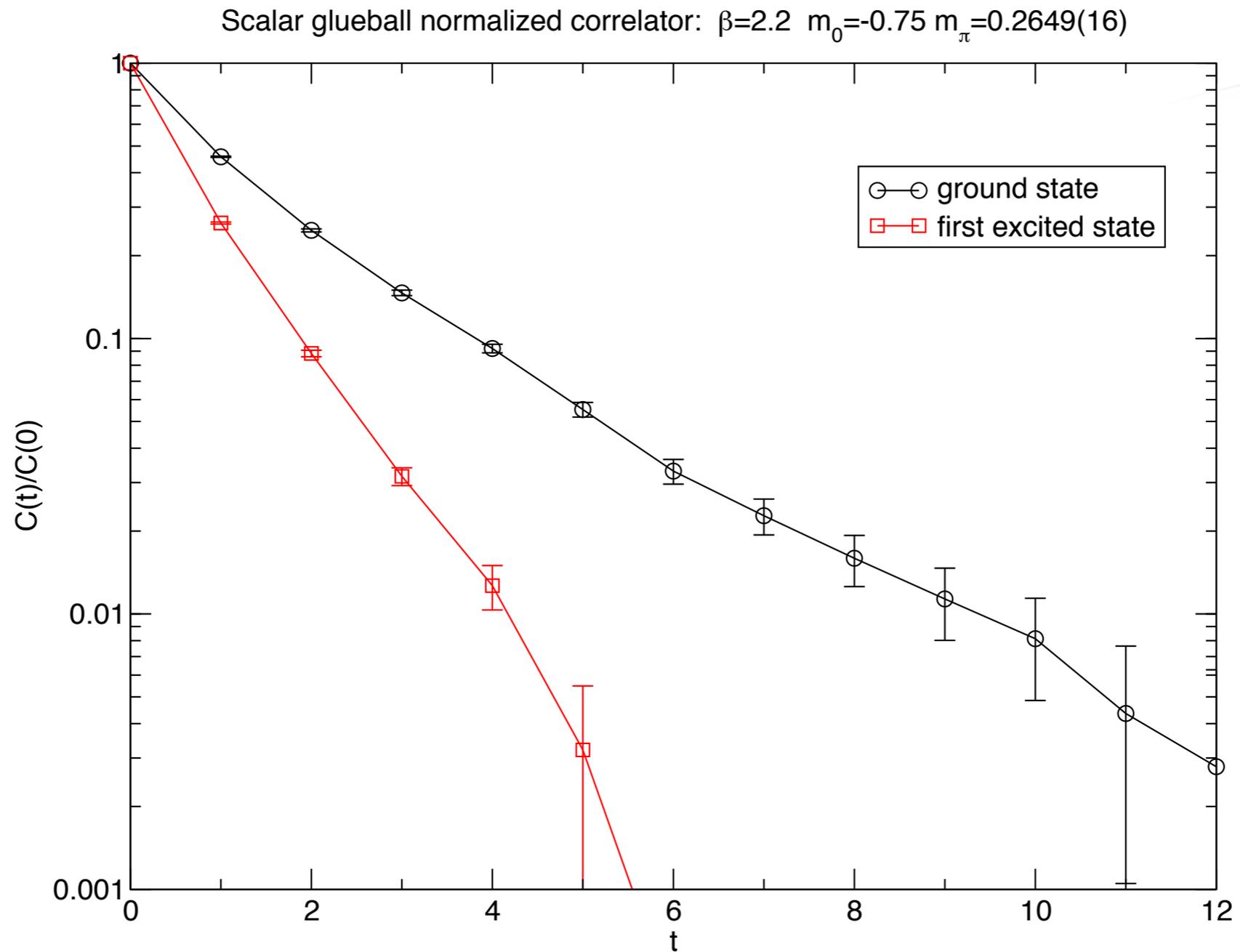
$$\Phi_\alpha(t) = \sum_{i=1}^n v_i^\alpha \mathcal{O}_i(t)$$

$$\langle \Phi_\alpha^\dagger(t) \Phi_\alpha(0) \rangle = |c_\alpha|^2 \left(e^{-m_\alpha t} + e^{-m_\alpha(T-t)} \right)$$

- the effective mass plateau is used to determine the fitting window on the correlator

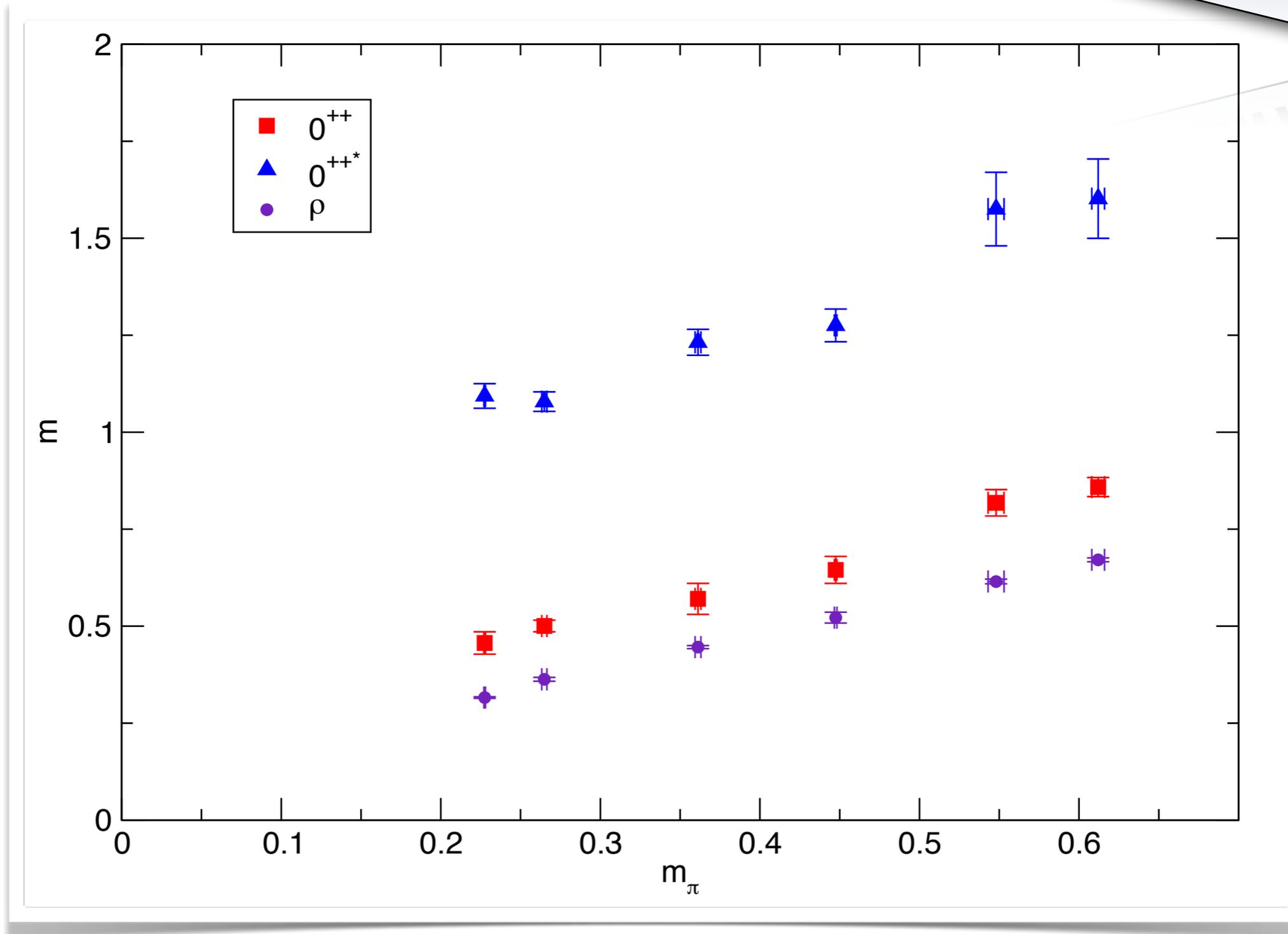
The scalar glueball

Preliminary



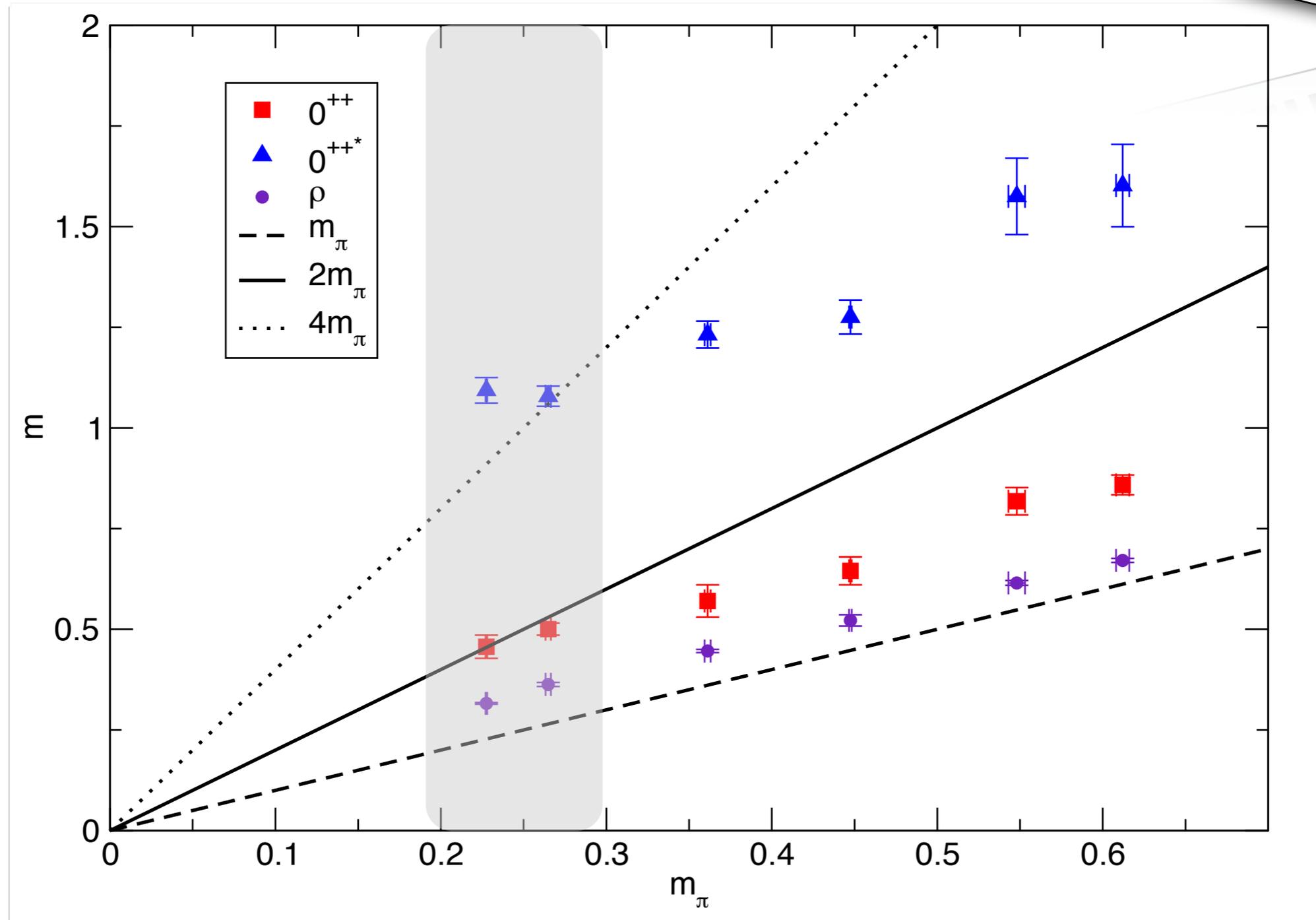
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Conclusions and future directions

- States in the **isosinglet scalar** channel have been investigated with **gluonic operators**, those coupling to glueballs in a pure gauge theory
- The ground and the first excited state are **heavier** than the isotriplet vector meson in the mass region explored
- Contributions due to **multi-pion** states become relevant in the light mass region
- Finite volume effects are being investigated for the lightest mass point on a $48^3 \times 96$ lattice
- A fermionic isosinglet scalar correlator, including the disconnected diagram, is being measured to address the mixing with the ground state

Thanks for your attention